

Problem 24:

$$1) Z(\mu', T) = \int_{\mathcal{R}} e^{-\beta \Omega + \beta \mu' \bar{N}}$$

$$= \sum_{\{n_{\vec{r}, i}\}} \prod_{\vec{r}} e^{-\beta \epsilon_{\vec{r}} n_{\vec{r}, 0} + \beta \mu' n_{\vec{r}, 0}} e^{-\beta (\epsilon_{\vec{r}} + \Delta) n_{\vec{r}, 1} + \beta \mu' n_{\vec{r}, 1}}$$

(over all configurations)

$$= \prod_{\vec{r}} \left(\sum_{n=0}^{\infty} e^{-\beta \epsilon_{\vec{r}} n + \beta \mu' n} \right) \left(\sum_{n=0}^{\infty} e^{-\beta (\epsilon_{\vec{r}} + \Delta) n + \beta \mu' n} \right)$$

$$\frac{1}{1 - e^{-\beta (\epsilon_{\vec{r}} - \mu')}} \quad \frac{1}{1 - e^{-\beta (\epsilon_{\vec{r}} + \Delta - \mu')}}$$

$$\Rightarrow \Omega = -k_B T \sum_{\vec{r}} \ln(1 - e^{-\beta (\epsilon_{\vec{r}} - \mu')}) - k_B T \sum_{\vec{r}} \ln(1 - e^{-\beta (\epsilon_{\vec{r}} + \Delta - \mu')})$$

$$= \Omega_0(\mu', T) + \Omega_0(\mu' + \Delta, T) \quad (2)$$

if $\Omega_0(\mu', T)$ is grand potential of "conventional" ideal Bose-Einstein gas

$$\text{for } \Omega_0: \langle n \rangle = \frac{1}{V} \frac{z}{1-z} + \frac{1}{\lambda_T^3} g_{3/2}(z) \quad z = e^{\beta \mu'}$$

$$\text{here: } \langle n \rangle = \frac{1}{V} \frac{z}{1-z} + \frac{1}{\lambda_T^3} g_{3/2}(z) + \frac{1}{V} \frac{z'}{1-z'} + \frac{1}{\lambda_T^3} g_{3/2}(z') \quad z' = e^{\beta (\mu' - \Delta)}$$

$$= \frac{1}{V} \frac{z}{1-z} + \frac{1}{\lambda_T^3} g_{3/2}(z) + \frac{1}{V} \frac{z e^{-\beta \Delta}}{1 - z e^{-\beta \Delta}} + \frac{1}{\lambda_T^3} g_{3/2}(z e^{-\beta \Delta})$$

$\xrightarrow{V \rightarrow \infty}$ even for $z=1$ (2)

maximal density in excited states at $z = z_{\text{max}} = 1$

$$\langle n \rangle_{\text{max}} = \frac{1}{\lambda_T^3} g_{3/2}(1) + \frac{1}{\lambda_T^3} g_{3/2}(e^{-\beta \Delta})$$

$$\Rightarrow \text{transition at } n_c = \frac{1}{\lambda_{T_c}^3} g_{3/2}(1) + \frac{1}{\lambda_{T_c}^3} g_{3/2}(e^{-\beta \Delta}) \quad (1)$$

if $\gamma \Delta \gg 1 \Rightarrow e^{-\gamma \Delta} \ll 1 \Rightarrow g_{3/2}(e^{-\gamma \Delta}) \approx e^{-\gamma \Delta}$

$$\lambda_{T_c}^3 n = g_{3/2}(1) + e^{-\frac{A}{k_B T_c}} \quad (1) \quad (*)$$

$$n \left(\frac{h^2}{2\pi m k_B T_c} \right)^{3/2} \leftarrow \text{equation for } T_c$$

2) ~~At $\Delta = 0$~~ Without the second state, we would have

$$\lambda_{T_c}^3 n = g_{3/2}(1)$$

With the second state the right hand side of (*) is larger than without the second state

$\Rightarrow T_c$ is lower

(2)

Problem 25:

a) one energy per \vec{p} , $\epsilon = \frac{\hbar^2}{2m} \left(\frac{2\pi}{L}\right)^2 \vec{p}^2$

→ number of \vec{p} 's below ϵ_F is

$$\frac{4}{3} \pi \left(\frac{2mL^2 \epsilon_F}{\hbar^2} \right)^{3/2}$$

with spin = $\frac{1}{2}$ this has to equal $\frac{N}{2}$

$$\frac{N}{2} = \frac{4}{3} \pi \left(\frac{2m \epsilon_F}{\hbar^2} \right)^{3/2} V$$

$$\frac{\hbar^2}{2m} \left(\frac{3}{8\pi} \frac{N}{V} \right)^{2/3} = \epsilon_F \quad (2)$$

$$\frac{N}{V} = 9 \frac{g}{\text{cm}^3} \cdot 63.5 \frac{g}{\text{mol}} \cdot 6 \cdot 10^{23} \frac{1}{\text{mol}} = 8.5 \cdot 10^{28} \frac{1}{\text{m}^3} \quad (2)$$

$$\epsilon_F = 1.13 \cdot 10^{-18} \text{ J} = 7 \text{ eV} \quad (1)$$

b) ~~$T_F = \frac{\epsilon_F}{k_B} = \frac{1.13 \cdot 10^{-18} \text{ J}}{1.38 \cdot 10^{-23} \text{ J/K}} = 8.2 \cdot 10^4 \text{ K}$~~

$$T_F = \frac{1.13 \cdot 10^{-18} \text{ J}}{k_B} = 8.2 \cdot 10^4 \text{ K} \quad (1)$$

Problem 26:

$$\begin{aligned} a) \quad \Omega(\epsilon) &= \# \text{ states with energy less or equal } \epsilon \\ &= \# \text{ } l\text{'s with } \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 l^2 \leq \epsilon \\ &= \pi \left(\frac{2m\epsilon L^2}{\hbar^2 (2\pi^2)} \right) = \frac{2\pi m}{\hbar^2} A \epsilon \end{aligned}$$

$$\Rightarrow n(\epsilon) = \frac{2\pi m}{\hbar^2} A \quad (2)$$

$$\begin{aligned} b) \quad N &= 2 \int_0^{\epsilon_F} n(\epsilon) d\epsilon = \frac{4\pi m}{\hbar^2} A \epsilon_F \Rightarrow \epsilon_F = \frac{2\hbar^2}{4\pi m} \frac{N}{A} = \frac{\hbar^2}{2\pi m} n \quad (1) \\ n(\epsilon) &= \frac{N}{2\epsilon_F} \quad (1) \end{aligned}$$

$$\begin{aligned} c) \quad N &= \int_{-\infty}^{\infty} n(\epsilon) \frac{2}{e^{\beta(\epsilon-\mu')} + 1} d\epsilon \quad (1) \\ &= \int_{-\infty}^{\infty} \Omega(\epsilon) \frac{2e^{\beta(\epsilon-\mu')}}{(e^{\beta(\epsilon-\mu')} + 1)^2} d\epsilon \\ &= \int_{-\infty}^{\infty} \Omega(\mu' + k_B T t) \frac{2e^t}{(e^t + 1)^2} dt \quad (1) \\ &= \int_{-\infty}^{\infty} \frac{N}{2\epsilon_F} \frac{\mu' + k_B T t}{\epsilon_F} \frac{2e^t}{(e^t + 1)^2} dt \\ &= \frac{N}{2} \cdot 2 \frac{\mu'}{\epsilon_F} \Rightarrow \mu' = \epsilon_F \quad (1) \end{aligned}$$

$$d) U = \int_{-\infty}^{\infty} n(\epsilon) \epsilon \frac{2}{e^{\beta(\epsilon - \mu')} + 1} d\epsilon \quad (1)$$

$$= \frac{N}{2\epsilon_F} \cdot 2 \int_{-\infty}^{\infty} \epsilon \frac{1}{e^{\beta(\epsilon - \mu')} + 1} d\epsilon$$

$$= \frac{N}{\epsilon_F} \int_{-\infty}^{\infty} \frac{\epsilon^2}{2} \frac{e^{\beta(\epsilon - \mu')}}{(e^{\beta(\epsilon - \mu')} + 1)^2} d\epsilon$$

$$= \frac{N}{2\epsilon_F} \int_{-\infty}^{\infty} (\mu' + k_B T t)^2 \frac{e^t}{(e^t + 1)^2} dt$$

$$= \frac{N}{2\epsilon_F} \epsilon_F^2 + \frac{N}{2\epsilon_F} (k_B T)^2 \frac{\pi^2}{3}$$

$$= \frac{N}{2} \epsilon_F + \frac{N\pi^2}{6} \epsilon_F \left(\frac{k_B T}{\epsilon_F} \right)^2 \quad (1)$$