

Closed system in contact with heat bath at temperature $T \Rightarrow$ average energy fixed

J11/15

Program: Maximize

$$\frac{S}{k_B} = - \int d\vec{x}^N g(\vec{x}^N) \ln[C_N g(\vec{x}^N)]$$

under the constraints

$$\int d\vec{x}^N g(\vec{x}^N) = 1$$

and $\int d\vec{x}^N H(\vec{x}^N) g(\vec{x}^N) = U$

J11/13

Method: Lagrange multipliers λ and μ

$$0 = \frac{\partial}{\partial g(\vec{x}^N)} \left[- \int d\vec{x}^N g(\vec{x}^N) \ln[C_N g(\vec{x}^N)] + \lambda \left(\int d\vec{x}^N g(\vec{x}^N) - 1 \right) + \mu \left(U - \int d\vec{x}^N H(\vec{x}^N) g(\vec{x}^N) \right) \right]$$

$$= - \ln[C_N g(\vec{x}^N)] - 1 + \lambda - \mu H(\vec{x}^N)$$

$$\Rightarrow g(\vec{x}^N) = \frac{1}{C_N} e^{\lambda-1} e^{-\mu H(\vec{x}^N)}$$

Normalization $\Rightarrow g(\vec{x}^N) = \frac{e^{-\mu H(\vec{x}^N)}}{\int d\vec{x}^N e^{-\mu H(\vec{x}^N)}}$

Still have to find μ :

$\mu = \mu(U) : U = \frac{\int d\vec{x}^N H(\vec{x}^N) e^{-\mu H(\vec{x}^N)}}{\int d\vec{x}^N e^{-\mu H(\vec{x}^N)}}$

$$\begin{aligned}
S &= -k_B \int d\vec{x}^N g(\vec{x}^N) \ln [C_N g(\vec{x}^N)] \\
&= -k_B \int d\vec{x}^N \frac{e^{-\mu H(\vec{x}^N)}}{\int d\vec{y}^N e^{-\mu H(\vec{y}^N)}} \left[\ln C_N + \mu H(\vec{x}^N) - \ln \int d\vec{y}^N e^{-\mu H(\vec{y}^N)} \right] \\
&= -k_B \ln C_N + k_B \mu U + k_B \ln \int d\vec{y}^N e^{-\mu H(\vec{y}^N)} \\
\frac{1}{T} = \frac{\partial S}{\partial U} &= k_B \mu + k_B U \frac{\partial \mu}{\partial U} = k_B \underbrace{\frac{\int d\vec{y}^N e^{-\mu H(\vec{y}^N)} H(\vec{y}^N)}{\int d\vec{y}^N e^{-\mu H(\vec{y}^N)}}}_{U} \frac{\partial \mu}{\partial U} \\
&= k_B \mu
\end{aligned}$$

$$\Rightarrow \mu = \frac{1}{k_B T} = \beta$$

Result:

$$g(\vec{x}^N) = \frac{e^{-\beta H(\vec{x}^N)}}{\int d\vec{y}^N e^{-\beta H(\vec{y}^N)}}$$

Abbréviation:

$$Z(T) = \int d\vec{x}^N \frac{e^{-\beta H(\vec{x}^N)}}{C_N}$$

partition function (dimensionless)

$$\begin{aligned}
S &= -k_B \ln C_N + k_B \frac{1}{k_B T} U + k_B \ln \int d\vec{x}^N e^{-\mu H(\vec{x}^N)} \\
&= \frac{U}{T} + k_B \ln Z(T)
\end{aligned}$$

$$\Rightarrow F = U - TS = -k_B T \ln Z(T)$$

→ same procedure as for discrete system

Useful property:

$$\begin{aligned}
U = F + TS &= F - T \left(\frac{\partial F}{\partial T} \right) = -k_B T \ln Z(T) - T \left(-k_B \ln Z(T) - k_B T \frac{\partial \ln Z(T)}{\partial T} \right) \\
&= k_B T^2 \left(\frac{\partial \ln Z(T)}{\partial T} \right) = - \left(\frac{\partial \ln Z(T)}{\partial \beta} \right)
\end{aligned}$$

V.1.4 Example (continuous system: ideal gas)

N particles, positions \vec{r}_i , momenta \vec{p}_i

$$H(\vec{x}^N) = \begin{cases} \frac{1}{2m} \sum_{i=1}^N \vec{p}_i^2 & \text{all } \vec{r}_i \text{ within volume } V \\ \infty & \text{otherwise} \end{cases}$$

$$Z(T) = \int d\vec{x}^N \frac{e^{-\beta H(\vec{x}^N)}}{C_N}$$

$$= \frac{1}{h^{3N} N!} \int d\vec{r}_1 \dots \int d\vec{r}_N \int d\vec{p}_1 \dots \int d\vec{p}_N e^{-\beta H(\vec{r}_1, \dots, \vec{r}_N, \vec{p}_1, \dots, \vec{p}_N)}$$

$$= \frac{1}{h^{3N} N!} V^N \int d\vec{p}_1 \dots \int d\vec{p}_N e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2}$$

$$= \frac{V^N}{h^{3N} N!} \left(\int d\vec{p} e^{-\frac{\beta}{2m} \vec{p}^2} \right)^N = \frac{V^N}{h^{3N} N!} \left(\frac{2\pi m}{\beta} \right)^{\frac{3N}{2}}$$

$$\Rightarrow F = -k_B T \ln Z(T) = -k_B T \ln \left[\frac{V^N}{h^{3N} N!} \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3N}{2}} \right]$$

$$= -N k_B T \ln \left[V \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} \right] + k_B T \ln N!$$

$$\approx -N k_B T - N k_B T \ln \left[\frac{V}{N} \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} \right] \quad \lambda \equiv \sqrt{\frac{h^2}{2\pi m k_B T}}$$

Stirling \rightarrow

$$= -N k_B T - N k_B T \ln \left(\frac{V}{N \lambda^3} \right)$$

$$P = - \left(\frac{\partial F}{\partial V} \right)_{T, N} = \frac{N k_B T}{V} \Rightarrow PV = N k_B T$$

"thermal wavelength"
quantum mechanical
wavelength of a
particle with energy
 $k_B T$

$$S = - \left(\frac{\partial F}{\partial T} \right)_{V, N} = N k_B + N k_B \ln \left[\frac{V}{N} \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} \right] + \frac{3}{2} N k_B$$

$$U = F + TS = \frac{3}{2} N k_B T$$

microscopic derivation of velocity distribution:

$$g(\vec{x}^N) = \frac{\frac{1}{c_V} e^{-\beta H(\vec{x}^N)}}{Z(T)} = \frac{h^{3N} N! e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2}}{h^{3N} N! V^N (2\pi m k_B T)^{3N/2}} \quad \text{if all } \vec{r}_i \text{ within } V$$

probability density to find a particle with velocity \vec{v} :

$$P_{\vec{v}}(\vec{v}) = \int d\vec{x}^N g(\vec{x}^N) \delta\left(\frac{\vec{p}_1}{m} - \vec{v}\right)$$

$$P_{\vec{v}}(\vec{v}) = \frac{1}{V^N (2\pi m k_B T)^{3N/2}} \int d\vec{x}^N e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2} \delta\left(\frac{\vec{p}_1}{m} - \vec{v}\right)$$

$$= \frac{1}{(2\pi m k_B T)^{3N/2}} \int d\vec{r}_1 \dots \int d\vec{r}_N \int d\vec{p}_1 \dots \int d\vec{p}_N e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2} \delta\left(\frac{\vec{p}_1}{m} - \vec{v}\right)$$

$$= \frac{1}{(2\pi m k_B T)^{3N/2}} \int d\vec{p}_1 \dots \int d\vec{p}_N e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2} \delta\left(\frac{\vec{p}_1}{m} - \vec{v}\right)$$

$$= \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left\{-\frac{m\vec{v}^2}{2k_B T}\right\}$$

→ Maxwell distribution

V.1.5 Example (mixed system)

Ideal gas of N identical atoms in box of volume V .
each atom has spin $\frac{1}{2}$ and magnetic moment μ .

external magnetic field \vec{H}

State of the system: $(\vec{x}^N, s_1, \dots, s_N)$ $s_i = \pm 1$

Energy : $E(\vec{x}, s_1, \dots, s_N) = \frac{1}{2m} \sum_{i=1}^N \vec{p}_i^2 - \frac{1}{2} \mu H \sum_{i=1}^N s_i$

partition function:

$$\begin{aligned}
Z(T) &= \frac{1}{h^N} \int d\vec{x}^N \sum_{\{s_1, \dots, s_N\}} e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2 + \frac{\beta \mu H}{2} \sum_{i=1}^N s_i} \\
&= \frac{V^N}{N! h^{3N}} \int d\vec{p}_1 \dots \int d\vec{p}_N \sum_{s_1 = \pm 1} \dots \sum_{s_N = \pm 1} e^{-\frac{\beta}{2m} \sum_{i=1}^N \vec{p}_i^2 + \frac{\beta \mu H}{2} \sum_{i=1}^N s_i} \\
&= \frac{V^N}{N! h^{3N}} \left(\int d\vec{p} e^{-\frac{\beta}{2m} \vec{p}^2} \right)^N \left(e^{-\frac{\beta \mu H}{2}} + e^{\frac{\beta \mu H}{2}} \right)^N \\
&= \frac{V^N}{N!} \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3N}{2}} \left(2 \cosh\left(\frac{\beta \mu H}{2}\right) \right)^N
\end{aligned}$$

↓ 1/15

$$\ln Z(T) = N \left\{ 1 + \ln \left[\frac{2V}{N} \left(\frac{2\pi m k_B T}{h^2} \right)^{\frac{3}{2}} \right] + \ln \cosh \frac{\beta \mu H}{2} \right\}$$

internal energy:

$$\begin{aligned}
U &= - \frac{\partial \ln Z}{\partial \beta} = + \frac{3}{2} N \frac{1}{\beta} = N \frac{\sinh \frac{\beta \mu H}{2}}{\cosh \frac{\beta \mu H}{2}} \frac{\mu H}{2} \\
&= \frac{3}{2} N k_B T = \frac{N \mu H}{2} \tanh \frac{\beta \mu H}{2}
\end{aligned}$$