

## IV.2 Microcanonical ensemble for N-particle system

N particles, energy E, volume V  
 Find density  $g(\mathcal{Z}^N)$  in thermodynamic equilibrium

Energy of phase space point  $\mathcal{Z}^N$ :  $H(\mathcal{Z}^N)$

energy continuous  $\rightarrow$  have to allow finite energy range  
 $[E, E + \Delta E]$   
 $\uparrow$  very small

$$g(\mathcal{Z}^N) = 0 \quad \text{if } H(\mathcal{Z}^N) < E \text{ or } H(\mathcal{Z}^N) \geq E + \Delta E$$

Constraint  $\int_{E \leq H(\mathcal{Z}^N) \leq E + \Delta E} d\mathcal{Z}^N g(\mathcal{Z}^N) = 1$

Isolated, closed system  $\Rightarrow$  Entropy is maximized

Maximize  $-\int_{E \leq H(\mathcal{Z}^N) \leq E + \Delta E} d\mathcal{Z}^N g(\mathcal{Z}^N) \ln C_N g(\mathcal{Z}^N) = F[g]$

$$G[g] \equiv \int_{E \leq H(\mathcal{Z}^N) \leq E + \Delta E} d\mathcal{Z}^N g(\mathcal{Z}^N) - 1$$

$$0 = \frac{d}{dg(\mathcal{Z}^N)} (F + \lambda G) = \begin{cases} -\ln C_N g(\mathcal{Z}^N) - 1 + \lambda & \text{for } H(\mathcal{Z}^N) \in [E, E + \Delta E] \\ 0 & \text{otherwise} \end{cases}$$

$$\Rightarrow C_N g(\mathcal{Z}^N) = e^{\lambda - 1} \quad \text{for } H(\mathcal{Z}^N) \in [E, E + \Delta E]$$

$$\Rightarrow g(\mathcal{Z}^N) \text{ is constant on "energy shell" } [E, E + \Delta E]$$

Normalization:

$$g(\vec{x}^N) = \begin{cases} \frac{1}{\Omega_{\Delta E}(E, V, N)} & H(\vec{x}^N) \in [E, E + \Delta E] \\ 0 & \text{otherwise} \end{cases}$$

where  $\Omega_{\Delta E}(E, V, N) = \int_{E \leq H(\vec{x}^N) \leq E + \Delta E} d\vec{x}^N =$  "volume of energy shell"

as  $\Delta E \rightarrow 0$  rewrite

"surface area of energy shell"   
  $\downarrow$  "surface"

$$\Omega_{\Delta E}(E, V, N) \approx \Delta E \frac{d\Omega(E, V, N)}{dE} \equiv \Delta E \Sigma(E, V, N)$$

with  $\Omega(E, V, N) = \int_{H(\vec{x}^N) \leq E} d\vec{x}^N =$  "phase space volume with energy less than E"

What is the entropy?

$$\begin{aligned} S &= -k_B \int d\vec{x}^N g(\vec{x}^N) \ln g(\vec{x}^N) C_N \\ &= -k_B \int_{E \leq H(\vec{x}^N) \leq E + \Delta E} d\vec{x}^N \frac{1}{\Omega_{\Delta E}(E, V, N)} \ln \frac{C_N}{\Omega_{\Delta E}(E, V, N)} \\ &= k_B \ln \frac{\Omega_{\Delta E}(E, V, N)}{C_N} \int_{E \leq H(\vec{x}^N) \leq E + \Delta E} d\vec{x}^N \frac{1}{\Omega_{\Delta E}(E, V, N)} \\ &= k_B \ln \frac{\Omega_{\Delta E}(E, V, N)}{C_N} \end{aligned}$$

Compare with discrete system:

$$\rightarrow \frac{\Omega_{\Delta E}(E, V, N)}{C_N} \equiv N(E) \text{ "number of states with energy E"}$$

$$\Rightarrow C_N = \text{"phase space volume of one state"}$$

Simplification of entropy for large N:

$$\Omega(E, V, N) = \int_0^E \Sigma(E', V, N) dE' \leq \int_0^E \Sigma(E, V, N) dE' = E \Sigma(E, V, N) = \frac{E}{\Delta E} \Omega_{\Delta E}(E, V, N)$$

↑  
energy surfaces increase

$$\Rightarrow \Omega_{\Delta E}(E, V, N) \leq \Omega(E, V, N) \leq \frac{E}{\Delta E} \Omega_{\Delta E}(E, V, N)$$

$$\Rightarrow \underbrace{\ln \Omega_{\Delta E}(E, V, N)}_{\sim N} \leq \underbrace{\ln \Omega(E, V, N)}_{\sim N} \leq \underbrace{\ln \Omega_{\Delta E}(E, V, N)}_{\sim N} + \underbrace{\ln \left( \frac{E}{\Delta E} \right)}_{\sim \ln N}$$

⇒ as N becomes large  $\ln \Omega_{\Delta E}(E, V, N) \approx \ln \Omega(E, V, N)$

$$\Rightarrow \boxed{S = k_B \ln \frac{\Omega(E, V, N)}{C_N}} \quad \text{for large } N$$

### IV.3 Summary

concepts: microcanonical ensemble  
phase space volume

facts: all states with the same energy are equally probable  
 $S = k_B \ln N(E) \quad | \quad S = k_B \ln \frac{\Omega(E, V, N)}{C_N}$

tools: Lagrange multipliers  
combinatorics  
Stirling's formula  
functional derivatives  
combinatorics →  $N(E, V, N) \rightarrow S(E, V, N) \rightarrow T(E, V, N) \rightarrow$  thermodynamics

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